Recent electron microscope studies of crystalline Nylon-6 specimens prepared from single crystals of the monomer reveal that twinning is of frequent occurrence in the polymer crystals (Fischer, 1965). The twinning appears to involve rotation of successive layers, or groups of layers, of the type shown in Fig.3, by $55^{\circ}$ relative to one another. Such twinning is not observed in Nylon specimens prepared from the melts (Macchi, 1966). Although the detailed mechanism of the twinning is not known, it is not difficult to relate the $55^{\circ}$ twin angle of the polymer to the $55^{\circ}$ by which molecular chains in successive layers of the monomer are inclined to one another.

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# The Crystal Structure of Thiodiglycollic Acid* 

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(Received 13 June 1966 and in revised form 3 January 1967)
Thiodiglycollic acid, $\mathrm{S}\left(\mathrm{CH}_{2} \mathrm{COOH}\right)_{2}$, crystallizes in space group Pnam with 4 molecules in the unit cell of dimensions $a=5.050 \pm 0.005, b=6.701 \pm 0.003, c=17.742 \pm 0.004 \AA$. The structure has been solved from two-dimensional photographic data and has been refined by least squares with three-dimensional diffractometer data. The structure consists of parallel infinite chains of hydrogen bonded molecules. Hydrogen bonding takes place between carboxyl groups around the centre of symmetry; they are coplanar within experimental error.

## Introduction

Thiodiglycollic acid, $\mathrm{S}\left(\mathrm{CH}_{2} \mathrm{COOH}\right)_{2}$ (also known as methyl sulphide $\alpha, \alpha$-dicarboxylic acid or thiodiacetic acid), is of importance because of its use as a precipitating organic reagent for zirconium (Sant \& Sant, 1959). It is used for the detection of copper, lead, mercury and silver, and also for the estimation of cadmium.

A preliminary report of the structure has already been published (Roy, 1965).

## Experimental

Single crystals of thiodiglycollic acid were grown by slow evaporation of a concentrated aqueous solution

[^0]at room temperature. The crystals appeared as thin plates elongated in the [100] direction. The unit-cell dimensions determined by means of a General Electric Single Crystal Orienter are: $a=5 \cdot 050 \pm 0 \cdot 005, b=6 \cdot 701$ $\pm 0.003, c=17.742 \pm 0.004 \AA$.
Systematic absences as found on Weissenberg photographs about the $a$ and $b$ axes were consistent with the space groups Pnam and Pna2 ${ }_{1}$. The density as determined by flotation is $1.66 \mathrm{g.cm}^{-3}$, and the calculated density for four molecules per unit cell is $1.70 \mathrm{~g} . \mathrm{cm}^{-3}$.
The structure was solved from the two projections along [100] and [010]. The $0 k l$ and $h 0 l$ intensities were estimated visually from the integrated Weissenberg photographs. The three-dimensional data used for the refinement of the structure were collected with a General Electric diffractometer equipped with a scintillation counter using $\beta$-filtered $\mathrm{Cu} K \alpha$ radiation and discrimination.
Peak intensities were measured which were converted to integrated intensities in the following way. For a set of 30 reflexions of $\theta$ values, distributed over the accessible $\theta$ range, both peak intensities and inte-
grated intensities were measured. Plotting the ratio of integrated intensity to peak intensity versus $\theta$ resulted in a graph which was used for converting peak intensities to integrated intensities.
To decide between the two possible space groups Pnam and $P n a 2_{1}$, the $0 k l$ reflexions were subjected to the $N(z)$ test (Howells, Phillips \& Rogers, 1950). Though the results seemed to indicate centrosymmetry they were not altogether convincing. The choice of Pnam was confirmed by the successful solution and refinement of the structure in terms of this space group.

## Structure determination and refinement

Approximate $y$ and $z$ parameters for sulphur were derived from the 0 kl Patterson projection. The fact that sulphur was found to occupy a special position on a mirror plane perpendicular to $\mathbf{c}$, together with a large concentration of peaks along the $c$ axis supported our choice of space group.
Starting from the sulphur parameters the projection was solved by the heavy atom technique. Successive Fourier and difference Fourier calculations reduced the $R$ index to $20 \%$.
The $x$ coordinate of sulphur was found from the $h 0 l$ Patterson projection. Approximate $x$ coordinates of the remaining atoms (except hydrogen) were obtained from bond length and bond angle considerations. This projection was refined as described above to an $R$ index value of $27 \%$. The [010] electron density projection is shown in Fig. 1.

The coordinates thus obtained were used as the starting point of a three-dimensional least-squares refinement with the diffractometer intensities, using Busing \& Levy's (1962) program. One overall scale factor was used and each observation assigned unit weight. Four cycles of refinement with individual isotropic temperature factors reduced $R$ from 22.5 to $14.5 \%$. Another four cycles, this time with individual anisotropic temperature factors, brought the discrepancy factor down to $11 \cdot 2 \%$. The temperature factor used was of the form
$\exp \left[-\left(\beta_{11} h^{2}+\beta_{22} k^{2}+\beta_{33} l^{2}+2 \beta_{12} h k+2 \beta_{13} h l+2 \beta_{23} k l\right)\right]$.
The positions of the three hydrogen atoms in the asymmetric unit were now located from two-dimensional difference Fourier maps. Four more cycles of refinement were done with the hydrogen atoms included, the thermal parameters of which were kept fixed. This reduced $R$ to $10.0 \%$.

Inspection of the observed and calculated structure factors showed that three reflexions, 020,110 , and 111, suffered from extinction effects: these were omitted from further refinement. Another cycle of refinement resulted in an $R$ value of $7.0 \%$.

Since the crystal was of appreciable dimensions ( $0.061 \mathrm{~cm} \times 0.020 \mathrm{~cm} \times 0.015 \mathrm{~cm} ; \mu=44.4 \mathrm{~cm}^{-1}$ ) an absorption correction was applied (Wehe, Busing \& Levy, 1962). One more cycle of refinement with the corrected structure factor gave no appreciable change in the $R$ value. Also a correction for anomalous dispersion was made (International Tables for X-ray Crystallography,


Fig.1. Electron density projection on (010). Contours are at equal arbitrary intervals. Full lines are positive contours, broken lines zero.

Table 1. Final atomic coordinates with estimated standard deviations

|  |  |  |  |  |  |  |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: |
| Atom | $x / a$ | e.s.d. | $y / b$ | e.s.d. | $z / c$ | e.s.d. |
| C(1) | 0.2813 | 0.0014 | -0.0228 | 0.0013 | 0.1749 | 0.0003 |
| $\mathrm{C}(2)$ | 0.1508 | 0.0012 | -0.0088 | 0.0011 | 0.0991 | 0.0002 |
| $\mathrm{O}(1)$ | 0.2931 | 0.0008 | -0.0873 | 0.0008 | 0.0458 | 0.0001 |
| $\mathrm{O}(2)$ | -0.0641 | 0.0008 | 0.0691 | 0.0008 | 0.0896 | 0.0002 |
| S | 0.0566 | 0.0004 | 0.0374 | 0.0004 | 0.2500 | 0.0000 |
| $\mathrm{H}(1)$ | 0.2254 |  | -0.0692 |  | -0.0107 |  |
| $\mathrm{H}(2)$ | 0.3198 | 0.0147 | -0.1663 | 0.0112 | 0.1785 | 0.0034 |
| $\mathrm{H}(3)$ | 0.4057 | 0.0137 | 0.0742 | 0.0109 | 0.1722 | 0.0034 |

Table 2. Anisotropic thermal vibration parameters with estimated standard deviations

| Atom | $\beta_{11}$ | e.s.d. | $\beta_{22}$ | e.s.d. | $\beta_{33}$ | e.s.d. | $\beta_{12}$ | e.s.d. | $\beta_{13}$ | e.s.d. | $\beta_{23}$ | e.s.d. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| C(1) | 0.0320 | $0 \cdot 0023$ | $0 \cdot 0200$ | $0 \cdot 0022$ | 0.0011 | $0 \cdot 0001$ | $-0.0023$ | 0.0028 | $-0.0003$ | 0.0005 | $0 \cdot 0000$ | 0.0005 |
| C(2) | 0.0259 | $0 \cdot 0026$ | 0.0186 | $0 \cdot 0024$ | $0 \cdot 0012$ | $0 \cdot 0001$ | 0.0009 | $0 \cdot 0024$ | 0.0007 | $0 \cdot 0005$ | $0 \cdot 0002$ | $0 \cdot 0004$ |
| $\mathrm{O}(1)$ | 0.0322 | $0 \cdot 0020$ | 0.0370 | $0 \cdot 0020$ | $0 \cdot 0009$ | $0 \cdot 0001$ | 0.0069 | $0 \cdot 0018$ | 0.0010 | 0.0003 | -0.0007 | $0 \cdot 0003$ |
| $\mathrm{O}(2)$ | 0.0339 | $0 \cdot 0021$ | 0.0319 | $0 \cdot 0022$ | $0 \cdot 0011$ | $0 \cdot 0001$ | $0 \cdot 0080$ | 0.0019 | $-0.0006$ | $0 \cdot 0006$ | -0.0013 | $0 \cdot 0003$ |
| S | 0.0269 | $0 \cdot 0009$ | $0 \cdot 0192$ | $0 \cdot 0007$ | $0 \cdot 0010$ | $0 \cdot 0000$ | $0 \cdot 0019$ | $0 \cdot 0008$ |  |  |  |  |



Fig. 2. Bond distances and bond angles. Standard deviations are given in brackets.
1962). This reduced the discrepancy index to $6.8 \%$. The final positions of all the atoms and their thermal parameters together with estimated standard deviations are given in Tables 1 and 2 respectively. A comparison

Table 3. Observed and calculated structure factors

of the observed and calculated structure factors is given in Table 3. The standard deviations of all the atoms were estimated from the diagonal elements of the inverse matrix.

## Results and discussion

## Molecular structure

The molecule has mirror symmetry, sulphur being on the mirror plane. The bond distances and angles in the molecule with their standard deviations are shown in Fig.2. Standard errors in bond angles were calculated by the method of Darlow (1960). Bond lengths and angles were calculated without correction for anisotropic thermal motion.

The sulphur-carbon bond length is $1 \cdot 80 \AA$, which does not differ significantly from the single bond value of $1.82 \AA$ mentioned by Abrahams (1956). At the sulphur atom the valency angle is $96^{\circ}$. This is considerably smaller than the corresponding angle in similar non-ring compounds (Abrahams, 1956; Stam, 1962). If we assume a van der Waals radius of $1 \cdot 2 \AA$ for H , there is no contact between the H atoms on $\mathrm{C}(1)$ and $\mathrm{C}\left(1^{\prime}\right)$, the distances $\mathrm{H}(2)-\mathrm{H}\left(2^{\prime}\right)$ and $\mathrm{H}(3)-\mathrm{H}\left(3^{\prime}\right)$ being 2.76 and $2.54 \AA$ respectively.

The carbon-carbon single bond length is $1.50 \AA$, which is in agreement with the values recently obtained for carboxylic acids (Eichhorn \& MacGillavry, 1959; Nardelli, Fava \& Giraldi, 1962; Higgs \& Sass, 1963; Okaya, 1965 ; van Eijck, Kanters \& Kroon, 1965). Within limits of error the configuration about $\mathrm{C}(1)$ is tetrahedral.

As can be seen from Fig. 2, $\mathrm{C}=\mathrm{O}$ is $1.22 \AA$ and $\mathrm{C}-\mathrm{O}$ is $1.30 \AA$. The angle $\mathrm{O}-\mathrm{C}=\mathrm{O}$ is $124^{\circ}$. This is in good agreement with the most recently published X-ray measurements on the carboxylic acids, as mentioned earlier.

It is not certain to what degree the structure of the carboxyl group is dependent on the molecular environment and hydrogen bond system. It has been suggested by Vaughan \& Donohue (1952) that the relative contributions of the two resonance forms depend upon the two oxygen atoms in the crystal structure, i.e. upon the number and strength of the hydrogen bonds. But considering other substances it is evident that this group will have very nearly the same structure in all carboxylic acids regardless of environment. In some acids there is no direct hydrogen bonding between the carboxylic groups.

## Crystal structure

The projection of the structure along [100] is shown in Fig.3. The structure consists of parallel infinite chains of molecules as shown in the Figure. The chains are formed by hydrogen bonding between
pairs of carboxyl groups around the centre of symmetry. The $\mathrm{O} \cdots \mathrm{O}$ distance is $2 \cdot 66 \AA$.
The equations of the plane containing the atoms $\mathrm{C}(2), \mathrm{O}(1), \mathrm{O}(2)$ and that containing $\mathrm{C}(1), \mathrm{C}(2), \mathrm{O}(1)$, $\mathrm{O}(2)$, together with their centrosymmetrically related ones (across the center of symmetry 000 ), are given in Table 4. The deviations of the atoms from these

Table 4. Least-squares planes
Plane $A$ (two carboxyl groups connected by hydrogen bonds)

$$
0.4317 x+0.8872 y-0.1630 z=0
$$

 related one)

$$
0.4337 x+0.8870 y-0.1568 z=0
$$

Deviations from the planes
Plane $A \quad$ Plane $B$

|  | $-0.0262 \AA$ | $-0.0100 \AA$ |
| :--- | :---: | :---: |
| $\mathrm{C}(1)$ | -0.028 | -0.0047 |
| $\mathrm{C}(2)$ | -0.0138 | 0.0159 |
| $\mathrm{O}(1)$ | 0.0095 | -0.0165 |
| $\mathrm{O}(2)$ | 0.0104 |  |

(each atom has been given unit weight.)

Table 5. Analysis of anisotropic temperature factors


Fig. 3. Projection of the structure on (100) showing molecular arrangement and some short intra- and inter-molecular distances.
planes are also given in the Table 4. These are leastsquares planes obtained by the method suggested by Schomaker, Waser, Marsh \& Bergman (1959). The atoms are co-planar within experimental error.
The $\beta_{i j}$ were used to calculate the displacements along the principal axes. At the same time the principal axes of the atomic ellipsoids and their direction cosines were determined. These data seem compatible with a libration of the carboxyl groups about the carbon-carbon bonds. This is given in Table 5.

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# The Identification and Crystal Structure of L-5-Methoxycarbonyl-7-formyl-1,2,5,6-tetrahydro$3 H$-pyrrolo[1,2-a]azepin-3-one, $\mathrm{C}_{12} \mathrm{H}_{13} \mathrm{NO}_{4}$, a Photolysis Product 

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The irradiation of $N$-chloroacetyl- $p$ - $O$-methyl-L-tyrosine resulted in the production of an unusual rearrangement product, m.p. $225^{\circ}$. The methyl ester of this product, m.p. $145^{\circ}$, has been identified by an X-ray analysis of its crystal structure to be


It crystallizes in the orthorhombic system, space group $P 2_{1} 2_{1} 2_{1}$ with four molecules per unit cell. The cell parameters are $a=7 \cdot 17, b=10 \cdot 08, c=15.99 \AA$. A partial structure was obtained by determining phases directly from the structure factor magnitudes by the use of the symbolic addition procedure as applied to noncentrosymmetric crystals. The complete structure evolved by the use of the tangent formula in a recycling procedure.

## Introduction

In the study of the effects of irradiation on $N$-chloroacetyl derivatives of amino acids Yonemitsu, Cerutti \& Witkop (1966) found that $N$-chloroacetyl-L-tryptophan



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